

Problem of the coexistence of several non-Hermitian observables in \mathcal{PT} -symmetric quantum mechanics

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Ilhem Leghrib

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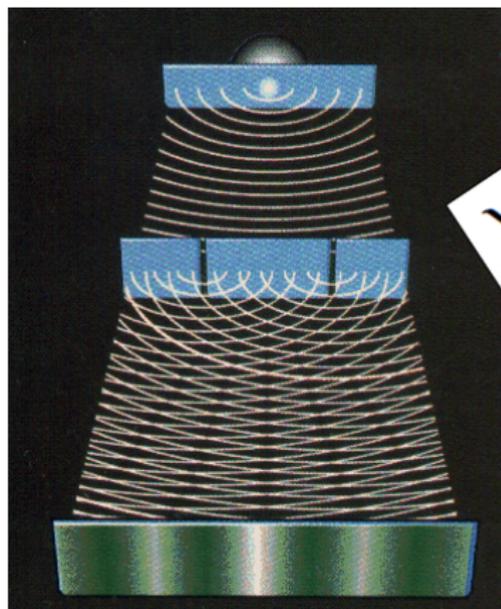
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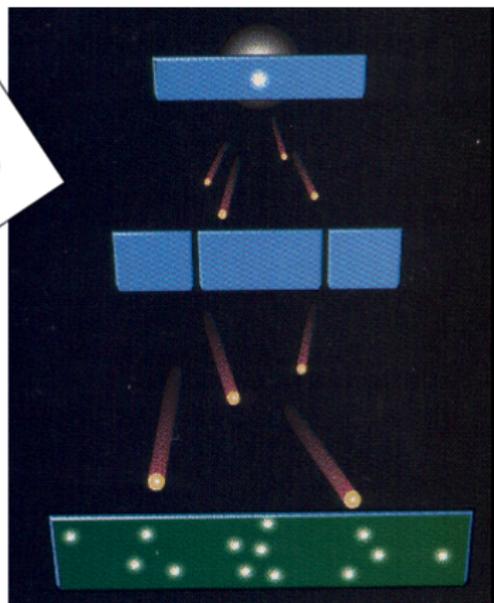
“Everything which is not forbidden is compulsory” Gell-Mann

- Basic Quantum Mechanics
- Some Terminology
- Background
 - PT-Symmetric Quantum Theory[1]
 - Complex Extension of Quantum Mechanics[2]
- Abstract of the paper
- Introduction of the paper
- Main result of the paper

Physical Motivation for Quantum Mechanics



$$\lambda = \frac{h}{p}$$



Mathematical Motivation for Quantum Mechanics

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- Can Hamiltonian be \mathcal{PT} -Symmetric?[1]

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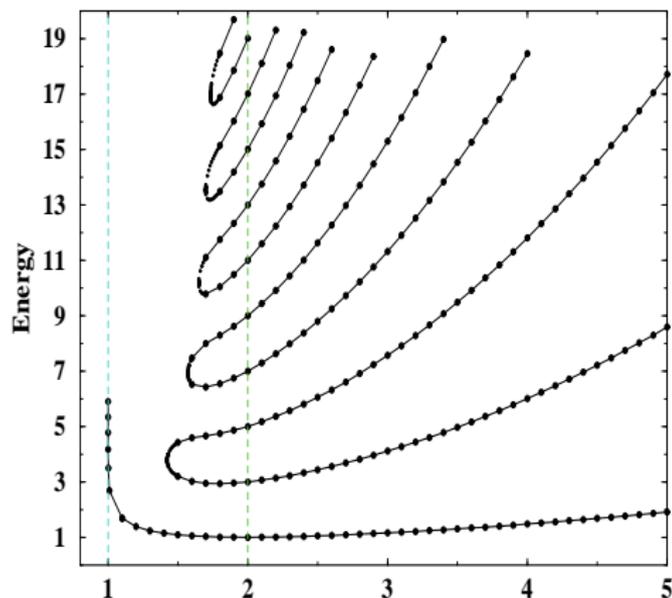
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- \mathcal{PT} -Symmetry is broken if $\delta < -2$



$$N = \delta + 2$$

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Now, $H\phi_n = E\phi_n \implies E\phi_n = E^*\phi_n$ [3] □

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- $\sum (-1)^n \phi_n(x)\phi_n(y) = \delta(x-y)$

\mathcal{PT} -Symmetry and Quantum Mechanics[1]

- $H = \frac{\hat{p}^2}{2m} + \frac{1}{2}k\hat{x}^2 (i\hat{x})^\delta$ has real eigenvalues for all $\delta \geq 0$ (note: H is \mathcal{PT} -symmetric but not Hermitian for $\delta \neq 0$)

Proof.

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- $\sum (-1)^n \phi_n(x)\phi_n(y) = \delta(x-y)$
- Can we now have a new condition: $H = H_{\mathcal{PT}}$ instead of $H = H^*$?

- If \mathcal{PT} -Symmetry is not broken, then is $\|H_{\mathcal{PT}}f\| = \|f\|$?

Complex Extension of Quantum Mechanics[2]

- If \mathcal{PT} -Symmetry is not broken, then is $\|H_{\mathcal{PT}}f\| = \|f\|$?
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- Choice: $\mathcal{C} = e^{Q(\hat{x}, \hat{p})} \mathcal{P}$ such that $Q(\hat{x}, \hat{p}) = -Q(-\hat{x}, -\hat{p})$ and $[\mathcal{C}, H] = 0$. Then, $\mathcal{C}^2 = 1$, $[\mathcal{C}, \mathcal{P}] \neq 0$ but $[\mathcal{C}, \mathcal{PT}] = 0$ so that $\|\phi\|_{\mathcal{CPT}} = 1$

During the recent developments of quantum theory it has been clarified that the observable quantities (like energy or position) may be represented by operators Λ (with real spectra) which are manifestly non-Hermitian in a preselected “friendly” Hilbert space $H^{(\mathcal{F})}$. The consistency of these models is known to require an upgrade of the inner product, i.e., mathematically speaking, a transition $H^{(\mathcal{F})} \rightarrow H^{(\mathcal{S})}$ to another, “standard” Hilbert space. We prove that whenever we are given more than one candidate for an observable (i.e., say, two operators Λ_0 and Λ_1) in advance, such an upgrade *need not* exist in general.

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- Natural restriction: $\Theta_j = U_j^* D_j U_j$ where U_j is an orthogonal matrix in the basis of the eigenvectors of an Λ_j
- When is $\Theta_j = \Theta_i$? Obvious: $[\Theta_j, \Theta_i] = 0$. Another possibility:
 $\Lambda_j \Theta_j = \Theta_j \Lambda_j^*$

Main Result

- Past approaches: let $\hat{x} \mapsto \hat{X}$, $\hat{p} \mapsto \hat{P}$ such that $\hat{a} = \frac{1}{\sqrt{2}} (\hat{x} + i\hat{p}) \mapsto \frac{1}{\sqrt{2}} (\hat{X} + i\hat{P})$ and $\hat{a}^\dagger = \frac{1}{\sqrt{2}} (\hat{x} - i\hat{p}) \mapsto \frac{1}{\sqrt{2}} (\hat{X} - i\hat{P})$ conditional to $\hat{a}\hat{a}^\dagger - q\hat{a}^\dagger\hat{a} = [\hat{a}, \hat{a}^\dagger]_q = I$. Let $q = 1 - \epsilon$.

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- $H_1 \neq H_1^*$ and $X_1 \neq X_1^*$

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